

$$i\hbar \frac{\partial}{\partial t} |\Psi\rangle = H |\Psi\rangle$$

$$\rightarrow i\hbar \frac{\partial}{\partial t} \Psi(x,t) = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \Psi(x,t) + V(x) \Psi(x,t) \quad \text{Seeking special solution}$$

remember  
light wave

$$\rightarrow e^{ikx} e^{-i\omega t} \quad \text{Solution: } |\Psi\rangle = |\phi_E\rangle c(t) \quad \leftarrow \text{complex}$$

$$i\hbar \frac{\partial}{\partial t} |\Psi\rangle = H |\Psi\rangle$$

$$i\hbar |\phi_E\rangle \frac{\partial}{\partial t} c(t) = \underbrace{H |\phi_E\rangle}_{E |\phi_E\rangle} c(t)$$

$$\Rightarrow i\hbar \frac{\partial}{\partial t} c(t) = E c(t) \rightarrow \frac{\partial}{\partial t} c(t) = -i \frac{E}{\hbar} c(t)$$

$$c(t) = e^{-i \frac{E}{\hbar} t}$$

$$-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \Psi(x,t) + V(x) \Psi(x,t) \Rightarrow \frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \phi_E(x) + V(x) \phi_E(x) = E \phi_E(x)$$

$$\Rightarrow \boxed{\Psi(x,t) = \phi_E(x) e^{-i \frac{E}{\hbar} t}} \quad \leftarrow \text{Stationary solution}$$

Example 1.)  $V(x) = 0, \quad \frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \phi_E(x) = E \phi_E(x)$

$$P: \phi_p(x) = e^{i \frac{p}{\hbar} x} \rightarrow p \phi_p(x)$$

$$\frac{\hbar}{i} \frac{\partial}{\partial x}$$

$$\frac{p^2}{2m} \phi_p(x) = \frac{p^2}{2m} \phi_p(x) = E \phi_p(x)$$

$$\Psi_p(x,t) = A \cdot e^{i \frac{p}{\hbar} x} e^{-i \frac{E}{\hbar} t} \quad \text{Solution for a free particle}$$

$$H \cdot \Psi_p(x,t) = \frac{p^2}{2m} \Psi_p(x,t) = E \Psi_p(x,t) = i\hbar \frac{\partial}{\partial t} \Psi_p(x,t) \quad \text{q.e.d.}$$

↙ expectation value (average) for measurement of  $\Omega$

Operator  $\Omega$ ,  $\langle \Omega \rangle = \langle \Psi | \Omega | \Psi \rangle$ , where  $\Psi$  could be  $\Psi(t)$

$$\langle \Psi(t) | \Omega | \Psi(t) \rangle = \int_{-\infty}^{\infty} \Psi^*(x,t) \Omega \Psi(x,t) dx$$

$$= \int_{-\infty}^{\infty} \phi_E^*(x) e^{i \frac{E}{\hbar} t} \Omega \phi_E(x) e^{-i \frac{E}{\hbar} t} dx$$

$$= e^{(i \frac{E}{\hbar} t - i \frac{E}{\hbar} t)} \int_{-\infty}^{\infty} \phi_E^*(x) \Omega \phi_E(x) dx$$

$$= \int_{-\infty}^{\infty} \phi_E^*(x) \Omega \phi_E(x) dx, \text{ no longer time-dependent } \Rightarrow \text{stationary}$$

GENERAL solution for free case:

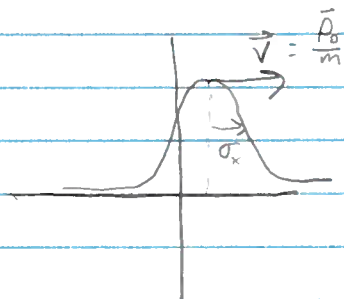
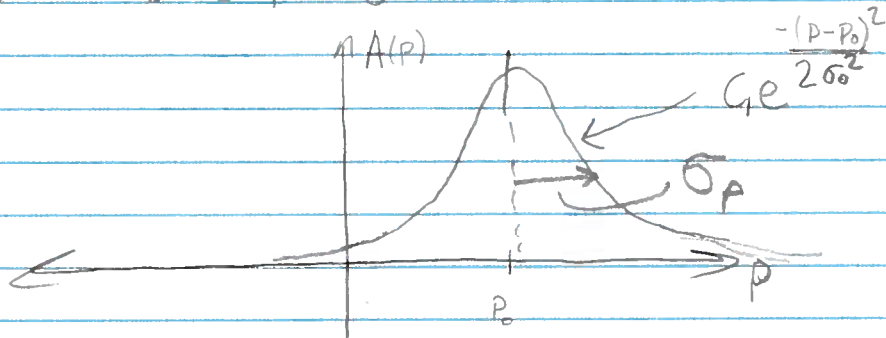
$$\Psi_{\text{free}}(x,t) = \int_{-\infty}^{\infty} dp A(p) e^{i \frac{p}{\hbar} x} e^{-i \frac{p^2}{2m\hbar} t}$$

← mixture of different plane waves (Fourier decomposition)

$$\begin{aligned} \text{In general: } H \phi_E(x) &= E \phi_E(x) \\ \Rightarrow \Psi(x,t) &= \int \phi_E(x) e^{-i E/\hbar t} \end{aligned}$$

Not an eigenstate to  $H$  (Hamiltonian operator)

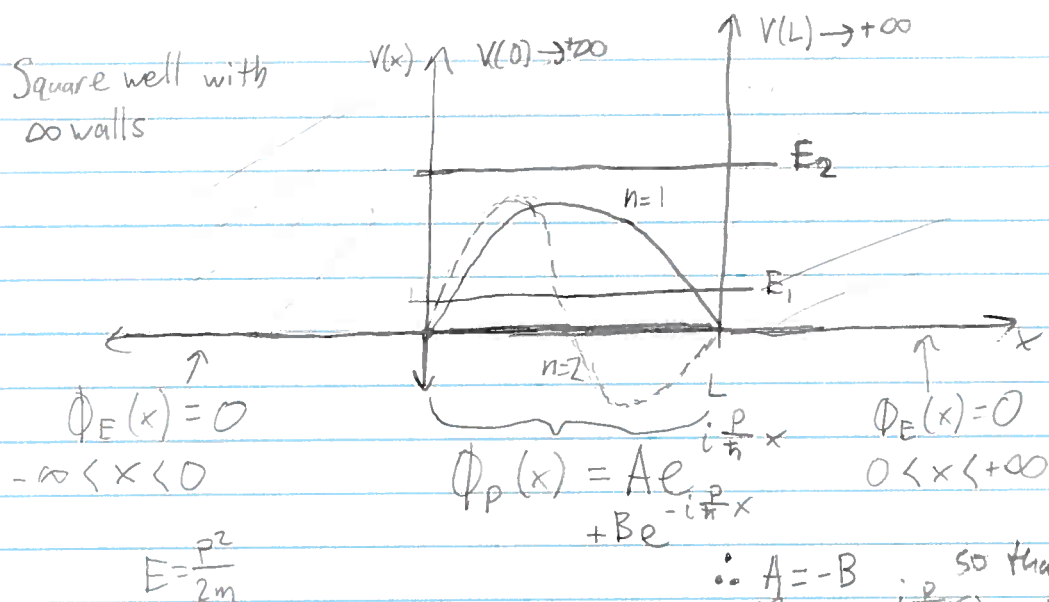
Ex: Gaussian wave package



For the Gaussian wave package

$$\sigma_x \cdot \sigma_p = \frac{\hbar}{2} \text{ for one point in time}$$

1-D Square well with  $\infty$  walls



$\therefore A = -B$  so that  $\Phi_p(0) = 0$   
 $\Phi_E(x) = A (e^{i \frac{p}{\hbar} x} - e^{-i \frac{p}{\hbar} x})$  [continuity!]

$\Phi_E^{(n)} = A \sin\left(\frac{p}{\hbar} x\right)$

$\sin\left(\frac{p}{\hbar} L\right) = 0$  [continuity @ L!]  
 $\frac{p}{\hbar} L = n\pi$  choose  $p_n = \frac{n\pi\hbar}{L}$

$H \Phi_{E_n}(x) = E_n \Phi_{E_n}(x)$

$E_n = \frac{n^2 \pi^2 \hbar^2}{2m L^2} = \frac{p_n^2}{2m}$

Note: there is a minimum  $E_1 = \frac{\pi^2 \hbar^2}{2m L^2}$ , not 0!

Energy is quantized:  $E_1, 4E_1, 9E_1, \dots$

